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449 one finds that for $y = F_0(e) C = :n; \sqrt{3} [P_0(2'Y_j) 3 -kjF(i) + (2'Y_j)! F_d(2'Y_j) 3 -ijF(-m,) 1 (14.17) C_2 = :n; [- (2'Y_j)! F_d(2'Y_j) 3 -ijF(i) + F_0(2'Y_j) 3 --; r(i)J$, and if y is to be $G_0(e)$, C and $G \sim (2'Y_j)$ replacing $P_0(2'Y_j)$ 1 2 and $G \sim (2'Y_j)$ replacing $F_d(2'Y_j)$. The values of the functions at $e_0 = 2'Y_j$ may be obtained from (14.8). 1 J.K. TYSON has employed the modified Hankel functions of order one third 2 as solutions of (13.4) to obtain expressions for the Coulomb functions for $L = 0$ which converge near $e = 2'Y_j$. His results appear as linear combinations of the real and imaginary parts of $n \sim(x) = (12)!e^{-x/6} [A;(-x) - iB;(-x)J$, (14.18) and its derivatives multiplying power series in $x = (e - 2'Y_j)j(2'Y_j)1$. For values 1 away from the turning point for $L = 0$, TYSON has obtained forms for $P_0(e)$ and $G_0(e)$ which are similar to (13.1) to (13.3). The JWKB approximation is again the leading term, and some higher order corrections are given. Expressions similar to Eqs. (14.11) and (14.12) have been obtained by T.D. 3 NEWTON employing the integral representation of (4.4). His results give re presentations of $FL(e)$, $Gde)$ in the vicinity of $e = 2'Y_j$ [whereas (14.11), (14.12) converge near $e = eLJ$ when L .

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