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Hamiltonian is.

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Homer Reid 's Solutions to Goldstein Problems: Chapter 8 2 From this we can immediately identify the T matrix and its inverse: $T = \begin{pmatrix} 2k & 2k \\ 2a + bq & 2l \end{pmatrix}$ $T^{-1} = \begin{pmatrix} 2k & 2k \\ 2a + bq & 2l \end{pmatrix}^{-1} = \frac{1}{(2k)^2 - (2a + bq)(2l)} \begin{pmatrix} 2l & -2k \\ -2k & 2a + bq \end{pmatrix}$ Then the Hamiltonian is $H = \frac{1}{2} (2a + bq)^2 \dot{\phi}^2 - k(2a + bq)^2 \phi + \frac{1}{2} p^2 + a + bq(2a + bq) - k(2a + bq) p + \frac{1}{2} p^2 - k(2a + bq) p + \frac{1}{2} p^2 - k(2a + bq) p + \frac{1}{2} p^2 - k(2a + bq) p$

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We use the first constraint to solve for the coordinate r : $r = R + a$, $\dot{r} = \dot{R}$, $\ddot{r} = \ddot{R}$. We use this solution in Lagrange 's equations for r , $\frac{\partial L}{\partial r} = -m(R + a) \ddot{r} + mg \sin \theta = m(R + a) \ddot{R} + mg(R + a) \cos \theta = \mu (R + a)$ (6) (7) We use the rolling constraint to find an expression for $\dot{\theta}$ as a function of \dot{R} : $\dot{\theta} = -\dot{R} + \dot{\phi}$ (8)

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Homer Reids Solutions to Goldstein Problems: Chapter 8 5 The kinetic energy is $m \dot{x}^2 + \frac{1}{2} m \dot{z}^2 = m (\dot{x} + L \cos \theta)^2 + \frac{1}{2} m \dot{z}^2$

$$+ L \sin)^2 2 m n o = (1 + 4a^2 x^2)x^2 + L^2 2 + 2Lx [\cos + 2ax \sin] .$$